tions, particularly of the type 00l, there existed the possibility that the observed characteristic extinctions were fortuitous. Therefore refinements were also performed in space groups P21 and *Pb.* In both cases the refinement converged to the same positional parameters and temperature factors as for *P21/b.* In addition the general set in both these non-centric space groups is twofold, two such sets being occupied by nickel and two by yttrium. The Busing, Martin & Levy program generates a correlation matrix during the refinement procedure. In both space groups this correlation matrix indicated correlation probabilities of 90-99% between twofold sets containing the same atomic species, and thus the correlation indicates a relationship which generates the fourfold sets of *P21/b.*

Discussion

The YNi structure is closely related to the orthorhombic ThNi structure (Florio, Baenziger & Rundle, 1956). The structural relationship can be shown by choosing a double unit cell for YNi with $a = 14.28$, $b=4.11$, and $c=5.50$ Å. Then in comparison the ThNi lattice parameters are $a=14.15$, $b=4.31$, and $c=$ 5.73 A. The atoms are arranged in this large YNi unit cell in two sets of the type $\pm (xyz; \frac{1}{4}+x, y, \frac{1}{2}-z;$ $\frac{1}{2}+x, y, z; \frac{3}{4}+x, y, \frac{1}{2}-z)$ with $x_{\text{N}i}=0.019, y_{\text{N}i}=0.245,$ $z_{\text{Ni}}=0.622$ and $x_{\text{Y}}=0.090$, $y_{\text{Y}}=0.249$, $z_{\text{Y}}=0.132$. Again in comparison the atoms in the ThNi structure are in two sets of the type $\pm (xyz; \frac{1}{4}+x, y, z; \frac{1}{2}+x, y, \frac{1}{2}-z;$ $\frac{3}{4} + x$, y , $\frac{1}{2} - z$) with $x_{\text{Ni}} = 0.018$, $y_{\text{Ni}} = 0.250$, $z_{\text{Ni}} = 0.630$ and $x_{\text{Th}}=0.094$, $y_{\text{Th}}=0.250$, $z_{\text{Th}}=0.140$. Thus a definite structural relationship is evident, the primary difference being in the sequence of the z parameters. A second significant difference is in symmetry with the lower symmetry of YNi resulting from the failure of the atoms to occupy the $y=\frac{1}{4}$ positions.

This failure of the yttrium and nickel atoms to occupy the $\frac{1}{4}$ and $\frac{3}{4}$ positions along the short axis is the only structural feature which differentiates the YNi structure from the B27 type structure (Pearson, 1958) typified by FeB (Bjurström & Arnfelt, 1929; Bjurström, 1933). Otherwise, the positional parameters and axial ratios of YNi are quite comparable to known B27 structures, and for classification purposes YNi should be included with this latter group.

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The Crystal Structure of the a-Modification of p-Nitrophenol near 90 °K

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The crystal structure of the α modification of p-nitrophenol near 90 °K has been determined by the use of partial three-dimensional data. The cell dimensions are $a = 11.66$, $b = 8.78$, $c = 6.098$ Å, $\beta = 107^{\circ}$ 32'. The space group is $P2_1/n$ and $Z = 4$. The crystals contain chains of hydrogen-bonded molecules. The benzene ring is planar but the nitrogen and oxygen atoms are displaced from the benzene-ring plane by amounts varying from $0.03-0.07$ Å.

Introduction

The structure analysis of the α modification of p-nitrophenol (I), started in 1957 and reported in this

paper, forms part of a research program on the chemistry of the solid state (Schmidt, 1957). The reaction occurring on irradiation of the α modification of p-nitrophenol manifests itself in the irreversible color change of the crystal from yellow to red; it is evidently structure dependent since the β -form is light stable as was first observed by Fritzsche (1859), and checked by us.

The room-temperature structure of the β modification will be described in a subsequent publication.

Experimental

Crystals of the α form can be grown by slow evaporation of solutions of p-nitrophenol in diethyl ether. The crystallographic constants near 300 and 90 °K are listed in Table 1. The cell dimensions compare well with previously reported room temperature values (Toussaint, 1954).

Table 1. *Crystallographic constants of x-p-nitrophenol*

$a = 11.8\text{\AA}$ $a = 11.66$ Å	
$b = 8.9$ $b = 8.78$ $c = 6.17$ $c = 6.098$	
$\beta = 106^{\circ} 52'$ $\beta = 107^{\circ} 32'$	
$P2\llap{1}$ P2, n $d_c = 1.51$ g.cm ⁻³ (Z = 4) $d_c = 1.551$ g.cm ⁻³ (Z = 4)	

* Cell dimensions determined from Weissenberg photographs, corrected for film shrinkage. Cell edges \pm 0·1%, β \pm 6′.

The $hk0$ and $hk3$ intensities were measured on a needle-shaped crystal of cross section 0.4×0.4 mm; for the *hOl* intensities we used a cube-shaped crystal. with an edge of about 0.3 mm, which had been cut by means of an alcohol-wetted nylon thread. All intensities were recorded at the temperature of boiling nitrogen with an equi-inclination Weissenberg camera (Hirshfeld & Schmidt, 1956) and Ni-filtered Cu K radiation. Intensities, measured visually, ranged from 28,000-1 for *h/cO,* 10,000-1 for *hO1,* and 2,000-1 for $hk3$. All usual corrections were applied by means of a computer program; no allowance was made for absorption. Approximate values of the scale factor k and of an average temperature factor were obtained from a Wilson plot (Wilson, 1942).

Solution of the structure

The $[hk0]$ zone was solved (Coppens, Schmidt & Gillis, 1958) by means of the inequality (Harker & Kasper, 1948)

$$
(U_{\mathbf{H}} \pm U_{\mathbf{H'}})^2 \leq (1 \pm U_{\mathbf{H} + \mathbf{H'}}). (1 \pm U_{\mathbf{H} - \mathbf{H'}}). \tag{1}
$$

The origin of the projection was fixed by the arbitrary assignment of signs to two reflections with h odd and k odd respectively. Systematic application of the inequality (1) by a procedure due to Gillis (1948) led to contradictions which could be climinatcd by a 20% reduction of the scale factor k. Subsequently, the signs of seven reflections were established absolutely; the signs of eleven other reflections could be expressed relative to each other. Two electron-density projections $\rho(xy)$ were computed with these twenty reflections; one of these maps could be interpreted and yielded trial coordinates whose correctness was established by the eventual refinement.

We may note here that unitary structure factors U were recalculated with the final scale and temperature factors. Somewhat surprisingly, their correct values are about half those used in the solution of the structure; the inequalities are now found to be nearly always satisfied by both signs, frequently with a *narrower* margin by the wrong sign *(e.g.* the example given below). Nevertheless, only one out of the eighteen signs determined with the old U 's proves to be incorrect. A similar experience has been recorded by Gillis (1948) who noted that the inequalities method had reserves of power in the sense that stronger inequalities than those used were in fact satisfied. He found that, if both signs satisfied an inequality, one by a comfortable margin and the other by a relatively narrow margin, the former sign was almost always the correct one.

An example of a sign determination is given below:

$$
|U_{410}| = 0.35; \ |U_{800}| = 0.70; \ |U_{020}| = 0.10
$$

$$
(U_{410} - U_{4\bar{1}0})^2 \le (1 - U_{800})(1 - U_{020}) \text{ since } U_{410} = -U_{4\bar{1}0}
$$

or $0.49 \leq (1-U_{800}) \times 1.1$ or $(1-U_{800}) \geq 0.45$ which is only satisfied if U_{800} is negative. (We have assumed U_{020} to be negative, thus choosing the less stringent condition. If U_{020} is taken as positive we get $(1-U_{800}) \ge 0.53$, which is a more powerful result.) The final values are however: $|U_{410}|=0.17$; $|U_{800}|=$

 $n =$ number of reflections; $s =$ number of parameters.

* These numbers are smaller than the sum of the corresponding values for the separate levels, since the reflections common to two levels were included only once in the refinement.

Table 3. Observed and calculated structure factors

 $Sign$ < indicates unobserved reflection

0.36; $|U_{020}| = 0.05$; now both signs of U_{800} satisfy the inequality, though with different margins.

According to the *(hk0)* projection the long molecular axis of the molecule near $(0, 0)$ lies roughly parallel to either [111] or $\overline{111}$. Approximate z coordinates for both orientations were derived from packing considerations and tested by structure-factor calculations; of the two models the former was found to correspond to the correct orientation.

Refinement

Details of our refinement procedure have been given in a previous publication (Coppens & Schmidt, 1964). The scattering factors employed were those given by Hoerni & Ibers (1954) for C, by Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955) for N and \overline{O} , and by McWeeny (1952) for H. Unobserved reflections were assigned a threshold value F_t based on the lowest intensity readings and only included in the refinement when $|F_c| > kF_t$. The three levels were first refined separately, initially with isotropic, later with anisotropic temperature parameters of all but the hydrogen atoms. When these refinements had been completed the levels were merged, and the residuals used to estimate the relative accuracy of the levels (Coppens & Schmidt, 1964). Numerical details are given in Table 2.

Refinement of all data was continued until the shifts in positional and thermal parameters were random and less than a fifth of the estimated standard deviations of these parameters. 360 reflections $(358$ observed reflections $+2$ unobserved reflections

Table 4. Final fractional coordinates of α -p-nitrophenol

Heavy atoms*

Hydrogen atoms*

* For the numbering of the heavy atoms see Figs. 2 and 3. $H(O)$ is the hydroxylic hydrogen; $H(n)$ is the hydrogen atom bonded to $C(n)$.

 \dagger Calculated assuming the aromatic HC= 1.08 Å, bisecting the angle between the adjacent C-C bonds, and the hydroxylie hydrogen at 1.00 Å from $O(3)$ towards the acceptor atom $O(2)$. with $|F_c| > kF_t$) were used to determine 113 parameters (3 scale factors, 45 positional and 65 thermal parameters).

Observed and calculated structure factors are listed in Table 3; final coordinates and temperature factors are given in Tables 4 and 5.

* The temperature factors β and β are respectively the coefficients in the expressions $\exp[-(\beta_{11}h^2+\beta_{22}k^2+\beta_{33}l^2+$ $\beta_{12}hk+\beta_{23}kl+\beta_{13}hl$] and exp [-(B sin² θ/λ^2)].

Fig. 1. Electron density projections. Contours each 1 e. \AA^{-2} , lowest contour at $2e^{\frac{1}{2}}$. (a) Projection along [001]. (b) Projection along [010].

Final overall agreement factors are

$$
\quad \text{and} \quad
$$

$$
r = \sum w(k^2 F_o^2 - F_c^2)^2 / \sum w k^4 F_o^4 = 3.5\%
$$

$$
R = \sum (|kF_o - |F_c||) / \sum kF_o = 8.3\%
$$

For 6 reflections $2|F_c| < kF_o$, while for 2 out of the total 35 unobserved reflections $F_c > kF_t$. Electron density projections and packing diagrams along [001] and

Fig. 2. Packing diagrams, (a) seen along [001], (b) seen along [010].

Discussion

Packing arrangement

The structure contains infinite chains of glide-plane related molecules linked together by hydrogen bonds. The angle between the benzene-ring planes of adjacent molecules in a chain is 74°.

Every molecule is faced by a second one at a distance of 3-33 A and related to it by a center of symmetry. Additional van der Waals contacts are made with twelve other molecules. Eleven of these are listed in Table 6; the two others are at x, y, z-1 and $\frac{1}{2}+x$, $\frac{1}{2} - y$, $z + \frac{1}{2}$. Distances to these two molecules are identical with those to molecules I and III given in Table 6. The molecule at $\bar{x}, \bar{y}, \bar{z}$ is at a somewhat larger distance, the shortest interatomic vector to it being $3.96~\text{\AA}$ (O(2)-O(2)).

Molecular structure

The benzene ring is planar, but the nitro group and the hydroxylic oxygen atom are displaced from the plane of the ring (Table 7). Such deviations from planarity are quite common in substituted aromatic molecules, as for example in benzoic acid (Sim,

Table 6. *Intermolecular distances* (Å)

Distances are given between reference molecule of which the atomic coordinates are listed in Table 4 and surrounding molecules

Distances involving hydrogen atoms are based on calculated hydrogen coordinates given in Table 4

* Hydrogen bond.

Table **7**

Equations of planes

Benzene-ring: $6.2340x - 5.2864y + 2.4667z - 0.2570 = 0$ Nitro-group: $6.4154x - 5.1074y + 2.4734z - 0.3358 = 0$

Angle between the two planes: $1^{\circ} 32'$

The benzene ring plane is derived according to Schomaker, Waser, Marsh & Bergman (1959)

Out-of-plane displacements

Robertson & Goodwin, 1955), p- dinitrobenzene (Abrahams, 1950), m-dinitrobenzene (Trotter, 1961) and nicotinic acid (Wright $\&$ King, 1953). Larger deviations from planarity are observed in the other modification of \overline{p} -nitrophenol (Coppens & Schmidt, 1965).

The intramolecular distances and angles are shown in Fig. 3. The average standard deviation in the bond lengths is 0.006 A, while the average standard deviation in the angles is 0.4° . The distances between the heavy atoms are all compatible with values recorded in the literature, and equal within the experimental errors with corresponding distances in the β form of p-nitrophenol. We shall discuss these dimensions in the paper on the structure of the β modification.

Fig. 3. Bond lengths and angles. Average standard deviation in bond lengths 0.006 Å, in angles 0.4° .

Hydrogen temperature factors

The fact that three out of five hydrogen 'temperature factors' are negative (Table 5) deserves some comment. The errors in these parameters are very large, but the negative sign indicates that the electron density is more concentrated in the bonded atom than in the free atom from which the f curve is derived. This is in agreement with recent calculations of the scattering factor of a hydrogen atom in a hydrogen molecule (Iijima & Bonham, 1963). Negative temperature factors for the hydrogen atoms have been observed in other low-temperature structures *(e.g.* Hirshfeld, Sandler & Schmidt, 1963).

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